

# Math-6500 Partial Differential Equations Fall 2007

## 1. PDEs of the first order

1.1. Write down an explicit solution for the initial-value problem

$$u_t + \mathbf{b} \cdot \nabla u + cu = 0, \quad x \in \mathbb{R}^n, t > 0, \quad u(x, 0) = g(x).$$

Here  $\mathbf{b}$  and  $c$  are constants and the gradient vector  $\nabla u = (\partial u / \partial x_1, \dots, \partial u / \partial x_n)$  is the gradient vector. **Hint:** As an easier particular case consider the transport equation, where  $n = 1$  and  $c = 0$ .

1.2. Solve the problem

$$xu_x + u_y = 1, \quad x \in \mathbb{R}, y > 0, \quad u(x, 0) = \exp(x)$$

1.3. Use the method of characteristics to solve the problem

$$xu_x + (x + y)u_y = 1,$$

where  $u(1, y) = y$  for  $0 < y < 1$ . Describe the region in the  $xy$ -plane over which the solution is uniquely determined.

1.4. Consider the differential equation in the *divergence form*:

$$\frac{\partial R(u)}{\partial y} + \frac{\partial S(u)}{\partial x} = 0, \quad (\text{div})$$

where  $R$  and  $S$  are known functions. [Note that the left-hand side is the divergence of the vector  $(S, R)$ .] Define a *weak solution*  $u(x, y)$  of (div) as a function for which the relation

$$\int \int (R(u)\phi_y + S(u)\phi_x) dx dy = 0 \quad (\text{integr})$$

holds for any function  $\phi(x, y)$  of class  $C_0^\infty$ , i.e. for any function  $\phi(x, y)$  that is infinitely differentiable and is zero everywhere except for a bounded set. (Relation (integr) follows formally from (div) through multiplication by  $\phi$  and integration by parts.)

Let  $u(x, y)$  be  $C^1$  in each of the two regions of the  $xy$ -plane separated by the curve  $x = \xi(y)$ . Show that if  $u(x, y)$  is a weak solution then:

- (a) It is a solution of (div) in each of the two regions;
- (b) It satisfies the following jump relation across the curve  $x = \xi(y)$  :

$$\frac{S(u^+) - S(u^-)}{R(u^+) - R(u^-)} = \frac{d\xi}{dy}$$

You may need to use the Stokes formula

$$\oint_{\partial B} P dx + Q dy = \int \int_B \left( \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dx dy$$

## 2. Properties of characteristics. General second-order equations.

**2.1.** Let  $u(x, t)$  be a  $C^1$ -solution of  $u_t + uu_x = 0$  in each of the two regions of the  $xy$ -plane separated by the curve  $x = \xi(t)$ . Let  $u$  be continuous but  $u_x$  have a jump discontinuity on the curve. Prove that then  $d\xi/dt = u$  and hence that the curve is a characteristic.

**Hint:** First prove that if  $u$  is continuous across the curve then the tangential derivative of  $u$ ,  $du(\xi(t), t)/dx$ , is also continuous. A similar argument also works for a more general equation  $a(x, y, u)u_x + b(x, y, u)u_y = c(x, y, u)$

**Comment:** A discontinuity in the derivative can be viewed as a disturbance in the otherwise smooth initial data; then the problem claims that the signal from this disturbance propagates along characteristics.

**2.2.** Let  $u(x, y)$  satisfy the equation

$$u_{xx} - 2u_{xy} + u_{yy} + 3u_x - u + 1 = 0$$

in a region of the  $xy$ -plane. Classify the equation and find its characteristics. Transform it into canonical form.

Construct a solution, if possible, for each of the following Cauchy problems:

(a)  $u = 2$ ,  $u_y = 0$  on the line  $y = 0$

(b)  $u = 2$ ,  $\partial u/\partial n = 0$  on the line  $y + x = 0$ , here  $\partial u/\partial n$  is the normal derivative, i.e. the directional derivative in the direction perpendicular to the line (find its expression via  $u_x$  and  $u_y$ ).

Are the data consistent with the governing equation in canonical form?

**2.3.** (Hodograph method or Legendre's transformation). Let  $u(x, y)$  be a solution of a quasi-linear equation of the form

$$a(u_x, u_y)u_{xx} + 2b(u_x, u_y)u_{xy} + c(u_x, u_y)u_{yy} = 0.$$

Introduce new independent variables

$$\xi = u_x(x, y), \quad \eta = u_y(x, y), \quad \phi = xu_x + yu_y - u.$$

Prove that  $\phi$  as a function of  $\xi$ ,  $\eta$  satisfies  $x = \phi_\xi$ ,  $y = \phi_\eta$  and the *linear* differential equation

$$a(\xi, \eta)\phi_{\xi\xi} - 2b(\xi, \eta)\phi_{\xi\eta} + c(\xi, \eta)\phi_{\eta\eta} = 0.$$

**Comment:** This trick is sometimes used to "linearize" nonlinear equations.

## 3. Fourier series

The problems in this section are pretty standard; they are taken from *Guenther & Lee*.

**3.1.** Problem 9, Sec. 3-1.

**3.2.** Problem 10, Sec. 3-1.

**3.3.** Problem 11, Sec. 3-1. Which properties of  $f(x)$  make one series better than the other?

**3.4.** Problem 7, Sec. 3-2.

#### 4. The wave equation in one spatial dimension

4.1. (A derivation of d'Alembert's solution). Show that the problem

$$u_{tt} - c^2 u_{xx} = 0, \quad u(x, 0) = f(x), \quad u_t(x, 0) = g(x)$$

can be represented as the system of two transport equations

$$u_t + cu_x = w, \quad w_t - cw_x = 0.$$

Find the initial conditions for the system that arise from the initial data for the wave equation; solve the system by the method of characteristics to obtain d'Alembert's formula.

4.2. (a) Solve the problem

$$u_{tt} - u_{xx} = 0, \quad u(x, 0) = f(x), \quad u_t(x, 0) = g(x),$$

when  $g(x) = 0$  and

$$f(x) = \begin{cases} x^2(1-x)^2, & 0 \leq x \leq 1 \\ 0, & \text{elsewhere} \end{cases}$$

Sketch the solution for  $t = 0, 1, 2$ .

(b) Repeat the problem in part (a) with different initial data:  $f(x) = 0$ ,

$$g(x) = \begin{cases} 1 - |x|, & |x| \leq 1 \\ 0, & \text{elsewhere} \end{cases}$$

4.3. Problem 6, Sec. 4-1 from *Guenther & Lee*. This problem explains how to split a solution of an inhomogeneous problem into a homogeneous and particular solutions (it is very similar to solving inhomogeneous ODEs). Problem 7 (not assigned!) explains, how to find a particular solution.

4.4. Problem 4, Sec. 4-2 from *Guenther & Lee* (a formal separation-of-variables Fourier series solution for the telegrapher's system). In addition, derive a second order equation for either  $i(x, t)$  or  $v(x, t)$ .

4.5. Problem 6, Sec. 4-3 from *Guenther & Lee*. (The idea is similar to the one in Problem 4.3).

4.6. Problem 2, Sec. 4-4 from *Guenther & Lee*. Use the methods of Sec. 4-4 to solve  $U_t + AU_x = 0$ , where  $U = (u_1, u_2)^{tr}$ ,

$$A = \begin{bmatrix} 4 & 1 \\ 3 & 2 \end{bmatrix}, \quad U(x, 0) = \begin{bmatrix} f(x) \\ 0 \end{bmatrix}, \quad f(x) = \begin{cases} (1-x^2)^2, & |x| \leq 1 \\ 0, & |x| > 1 \end{cases}$$

Discuss the behavior of the waves for  $t > 0$ .

4.7. Find a d'Alembert-type solution to the initial-boundary value problem

$$\begin{aligned} u_{tt} - c^2 u_{xx} &= 0, \quad x > 0, t > 0 \\ u(x, 0) &= f(x), \quad u_t(x, 0) = g(x), \quad x \geq 0, \\ u(0, t) &= h(t), \quad t \geq 0 \end{aligned}$$

by converting it to a first-order system and applying the results of Sec. 4-5 (*Guenther & Lee* Problem 5(a), Sec. 4-5).

## 5. Primer of distributions ( $\mathcal{D}'$ )

**5.1.** Let  $f(\xi)$  be a *continuous* (not necessarily differentiable) function. Show that  $u(x, t) = f(x - ct)$  is a distribution (weak) solution of the partial differential equation  $u_t + cu_x = 0$  (F. John, Problem 3.6.1). **Hint:** Transform the integral

$$\int \int (\phi_t + c\phi_x)u \, dx \, dt$$

to the coordinates  $y_1 = x - ct$ ,  $y_2 = t$ . It is sufficient to use special test functions  $\phi = \sigma(y_1)X(y_2)$ .

**5.2.** Prove that the function  $u(x_1, x_2)$  defined by

$$u(x_1, x_2) = \begin{cases} 1, & \text{for } x_1 > \xi_1, x_2 > \xi_2 \\ 0, & \text{elsewhere} \end{cases}$$

is a fundamental solution with the pole  $(\xi_1, \xi_2)$  of the operator  $L = \partial^2/\partial x_1\partial x_2$  in the  $x_1x_2$ -plane. (F. John, Problem 3.6.2)

**5.3.** Prove that the function  $u(x_1, x_2)$  defined by

$$u(x_1, x_2) = \begin{cases} 1/2, & \text{for } |x_1 - \xi_1| < \xi_2 - x_2 \\ 0, & \text{elsewhere} \end{cases}$$

is a fundamental solution with the pole  $(\xi_1, \xi_2)$  of the operator  $L = \partial^2/\partial x_2^2 - \partial^2/\partial x_1^2$  in the  $x_1x_2$ -plane. (F. John, Problem 3.6.3)

**5.4.** Prove that  $\lim_{n \rightarrow \infty} \sin(nx) = 0$  in the space of distributions  $\mathcal{D}'(\mathbb{R})$ , i.e., that

$$\lim_{n \rightarrow \infty} \int \sin(nx)\phi(x) \, dx = 0$$

for any  $\phi \in C_0^\infty$ . Find  $\lim_{n \rightarrow \infty} \sin^2(nx)$  in the space of distributions  $\mathcal{D}'(\mathbb{R})$ .

**Comment:** This example shows that multiplication of distributions is not continuous, even when it is defined.

**5.5. (a)** Prove that  $x\delta(x) = 0$  in the sense of distributions. Thus, in  $\mathcal{D}'(\mathbb{R})$  the equation  $xT(x) = 0$  has a solution  $T(x) = c\delta(x)$  (cf. *Guenther & Lee* Problem 10-5.4).

**(b)** Solve the differential equation  $dT/dx = \delta(x)$  (the solution then is a fundamental solution of the equation). Don't forget to prove that your solution satisfies the equation in  $\mathcal{D}'(\mathbb{R})$ .

## 6. The Heat Equation

**6.1.** (Duhamel's Principle and variation of parameters). Consider the initial value problem

$$\frac{dx}{dt} + A(t)x = f(t), \quad x(0) = 0, \quad x \in \mathbb{R}^n$$

Let  $\phi^{(1)}(t, \tau), \dots, \phi^{(n)}(t, \tau)$  (each  $\phi$  is a column vector) be a set of solutions of the homogeneous problem

$$\frac{dx}{dt} + A(t)x = 0,$$

with the initial conditions

$$\phi^{(1)}(t, \tau)|_{t=\tau} = (1, 0, \dots, 0)^{tr}, \dots, \phi^{(n)}(t, \tau)|_{t=\tau} = (0, 0, \dots, 1)^{tr}$$

Thus, the matrix  $\Phi(t, \tau) = [\phi^{(1)} \dots \phi^{(n)}]$  satisfies the (matrix) equation

$$\begin{aligned} \frac{d\Phi(t, \tau)}{dt} + A(t)\Phi(t, \tau) &= 0, \\ \Phi(t, \tau)|_{t=\tau} &= I \end{aligned}$$

- (a) (Semigroup property.) Show that if  $t_1 < t_2 < t_3$  then  $\Phi(t_3, t_1) = \Phi(t_3, t_2)\Phi(t_2, t_1)$ .  
 (b) Look for a solution of the inhomogeneous problem in the form

$$x = c_1(t)\phi^{(1)} + \dots + c_n(t)\phi^{(n)} = \Phi c$$

Obtain a differential equation for  $c(t)$  and integrate it to obtain a representation of the solution via  $x(t)$  as an integral. Interpret this integral as Duhamel's integral (cf. *Guenther & Lee* Sec.5.3).

**6.2.** *Guenther & Lee* Problem 3, Sec. 5-3. This an analogue of the semigroup property [problem 6.1(a)] for the heat equation.

The next two problems suggest a method to solve the following problem

$$\begin{aligned} u_t &= u_{xx}, \quad t > 0, \quad x > 0 \\ u(x, 0) &= 0, \quad u(0, t) = \mu(t) \end{aligned} \tag{1}$$

The result is another Duhamel's integral.

**6.3.** Assume that there exists a one-parametric family of functions  $U(x, t; \tau)$  ( $\tau$  is a parameter) that solve the heat equation

$$u_t = u_{xx}, \quad t > \tau, \quad x > 0$$

with the initial conditions at  $t = \tau$  and boundary conditions at  $x = 0, t > \tau$

$$u(x, \tau) = 0; \quad u(0, t) = 1, \quad t > \tau.$$

- (a) Use  $U$  to solve the problem

$$\begin{aligned} u_t &= u_{xx}, \quad t > t_1, \quad x > 0; \\ u(x, t_1) &= 0; \quad u(0, t) = 1, \quad t_1 < t < t_2, \quad u(0, t) = 0, \quad t > t_2. \end{aligned}$$

- (b) Use the result in part (a) to justify the formula for the solution of (1)

$$u(x, t) = \int_0^t \frac{\partial}{\partial \tau} U(x, t; \tau) d\tau$$

6.4. Use the similarity method to solve the problem

$$u_t = u_{xx}, \quad t > 0, \quad x > 0; \quad u(x, 0) = 0; \quad u(0, t) = 1, \quad t > 0.$$

Namely, look for a solution  $U(x, t) = v(x/\sqrt{t})$  and reduce the problem to a boundary value problem for an ordinary differential equation for  $v(\xi)$ , where  $\xi$  is the similarity variable. How does  $U(x, t; \tau)$  from problem 6.3 relate  $U(x, t)$ ?

6.5. Problem 5-4 #7 from *GL*.

6.6. Energy method for the heat equation. Consider the problem

$$u_t = u_{xx}, \quad t > 0, \quad L > x > 0, \quad u(x, 0) = f(x) \tag{2}$$

with either  $u(0, t) = u(L, t) = 0$  or  $u_x(0, t) = u_x(L, t) = 0$ . You may assume as much differentiability from  $u$  as you want.

(a) Show that the "energy"

$$E(t) = \int_0^L u^2(x, t) dx$$

decays as a function of  $t$ .

(b) Deduce from (a) uniqueness for the problem (2) with boundary conditions  $u(0, t) = f(t)$ ,  $u(L, t) = g(t)$ .

Comment: This energy has nothing to do with any physical quantity of any significance.

## 7. Integral Equations

7.1. Consider the integral equation

$$u(x) = f(x) + \lambda \int_0^{10} xtu(t)dt$$

(a) For which  $\lambda$  does the equation have a unique solution for any  $f$ ?

(b) For which  $\lambda$  does the equation have no solution for a generic  $f$  and infinitely many solutions for a "special" admissible  $f$ ? Describe the admissible functions  $f$  and solve the equation.

7.2. Volterra integral operators: Problem 7-4 #9(a-d) from *GL*.

7.3. Solve the integral equation

$$y(t) = t + \int_0^1 (1 + st)y(s)ds$$

## 8. The Laplace Equation

1. Ill-posedness of the Cauchy problem for Laplace's equation (*cf. Guenther & Lee Sec. 1-8*). In two dimensions consider the problem

$$\begin{aligned} \Delta u &= 0, \quad \infty < x < -\infty, \quad y > 0, \\ u(x, 0) &= f(x), \quad u_y(x, 0) = g(x), \quad \infty < x < -\infty \end{aligned}$$

Construct a sequence of separated solutions

$$u_n(x, y) = \frac{1}{n} Y_n(y) \cos nx$$

such that  $u_n(x, 0) \rightarrow 0$ ,  $(u_n)_y(x, 0) \rightarrow 0$  as  $n \rightarrow \infty$ , while  $u_n(x, 1) \rightarrow \infty$ . This example shows that the solution doesn't depend continuously on the initial data: indeed, the data converge

$$F_n(x) = f(x) + u_n(x, 0) \rightarrow f(x), \quad G_n(x) = g(x) + (u_n)_y(x, 0) \rightarrow g(x)$$

but the corresponding solutions do not converge  $u(x, y) + u_n(x, y) \not\rightarrow u(x, y)$ .

2. (Based on Problem 4.1 #2 of *Partial Differential Equations* by F. John). Let  $L = \Delta + c$  in  $n = 3$  dimensions, where  $c > 0$  is a constant. ( $L$  is Helmholtz or reduced wave operator).

(a) Find all solutions of  $Lu = 0$  with spherical symmetry. Hint: Set  $u(r) = v(r)/r$ .

(b) Prove that

$$\Phi(x) = -\frac{\cos(\sqrt{c}|x|)}{4\pi|x|}$$

is a fundamental solution for  $L$ . In other words,

1. Use the divergence theorem to prove that

$$\int_{|x|<a} L\Phi \, dx = 0$$

2. \* Repeat the argument from *Evans* (paragraphs 2–4, pages 24–25) to demonstrate that

$$u(x) = \int_{\mathbb{R}^3} \Phi(x-y)f(y) \, dy$$

solves  $Lu = f$ .

3. (Based on Problems 8-2 #5,6 of *Guenther & Lee*).

(a) In two dimensions use Fourier transform to solve formally

$$\Delta u = 0, \quad u(x, 0) = f(x), \quad \infty < x < -\infty, \quad y > 0,$$

and so find, assuming suitable behavior of  $u$  at infinity (state your assumption), that

$$u(x, y) = \frac{y}{\pi} \int \frac{f(s)ds}{(x-s)^2 + y^2}$$

for  $y > 0$ .

(b) Now show that the formal solution of part (a) solves the problem. Namely:

1. Check directly that the solution  $u(x, y)$  satisfies  $\Delta u = 0$ , for  $\infty < x < -\infty$ ,  $y > 0$ .

2. Prove that  $u(x, y) \rightarrow f(x)$  as  $y \rightarrow 0 +$ .

4. *Guenther & Lee* Problem 11, Sec. 8-3. The solution formula in (3-13) gives a solution of the Dirichlet problem for the Poisson equation through the dipole and volume potentials that uses Green's function  $G$  for this problem. Define a Neumann  $N(x, y)$  and Robin  $R(x, y)$  functions that provide similar representations for the Neumann and Robin problems.

5. (*John* 4.2 #1) Let  $\Omega$  denote the domain  $|x| > 1$ . Let  $u \in C^2(\overline{\Omega})$ ,  $\Delta u = 0$  in  $\Omega$ , and  $\lim_{x \rightarrow \infty} u(x) = 0$ . Show that

$$\max_{\overline{\Omega}} |u| = \max_{\partial\overline{\Omega}} |u|. \quad (3)$$

6. (Cf. *GL* 8-7, Problem 2) Consider the interior Robin problem with  $\alpha = \text{const}$ . Assume a solution in the form of a single-layer potential. Use the jump relations to find the integral equation for the density of the potential.